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Aalto University

Linear time-invariant processes: Dynamics CHEM-E7190 (was E7140), 2021

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Representation and analysis

Consider a linear and time-invariant system of order N_x , in state-space representation

 \rightsquigarrow Let N_x be the number of outputs \rightsquigarrow Let N_u be the number of inputs $\begin{cases} \dot{x}(t) = Ax(t) + Bu(t) \\ y(t) = Cx(t) + Du(t) \end{cases}$

System

A $(N_x \times N_x)$, B $(N_x \times N_u)$, C $(N_y \times N_x)$ and D $(N_y \times N_u)$ are the system matrices $\sim x(t)$ is the state vector

- $(N_x \text{ components})$
- $\rightsquigarrow \dot{x}(t)$ is the derivative of the state vector
- $(N_x \text{ components})$
- $\rightsquigarrow u(t)$ is the input vector
- $(N_u \text{ components})$
- $\rightsquigarrow y(t)$ is the **output vector**
- $(N_y \text{ components})$

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Representation and analysis (cont.)

The analysis problem: Determine the behaviour of state x(t) and output y(t) for $t \ge t_0$

- We are given the input function u(t), for $t \ge t_0$
- We are given the initial state $x(t_0)$

The solution to the analysis, for $t \ge t_0$, an initial state $x(t_0)$ and an input $u(t \ge t_0)$

$$x(t) = e^{A(t-t_0)}x(t_0) + \int_{t_0}^t e^{A(t-\tau)}Bu(\tau)d\tau$$
$$y(t) = \underbrace{Ce^{A(t-t_0)}x(t_0) + C\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)d\tau}_{Cx(t)} + Du(t)$$

The solution is known as the Lagrange formula

• Based on the state transition matrix

 $\rightsquigarrow e^{At}$

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Note that we can write the state solution x(t), for $t \ge t_0$, as the sum of two terms

$$x(t) = \underbrace{e^{A(t-t_0)}x(t_0)}_{x_u(t)} + \underbrace{\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)d\tau}_{x_f(t)}$$

= $x_u(t) + x_f(t)$

- \rightsquigarrow The force-free evolution of the state, $x_u(t)$
- \rightsquigarrow The forced evolution of the state, $x_f(t)$

The force-free evolution of the state, from the initial condition $x(t_0)$ $\rightsquigarrow e^{A(t-t_0)}$ determines the transition from $x(t_0)$ to x(t)

 \rightsquigarrow In the absence of contribution from the input

The forced evolution of the state, from the contribution of input u(t) \rightsquigarrow In the absence of an initial condition $x(t_0)$

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The state transition matrix

Consider a square $(N_x \times N_x)$ matrix A, the exponential e^A is square $(N_x \times N_x)$ matrix

$$e^A = I + A + \frac{A^2}{2!} + \frac{A^3}{3!} + \dots$$
$$= \sum_{k=0}^{\infty} \frac{A^k}{k!}$$

The state transition matrix is the matrix exponential e^{At} of the matrix At

 $\rightsquigarrow\,$ It is a matrix whose elements are functions of time

 \rightsquigarrow We discuss its meaning and how to compute it

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The state transition matrix (cont.)

The exponential function

Let z be some scalar, by definition its exponential is a scalar

$$e^{z} = 1 + z + \frac{z^{2}}{2!} + \frac{z^{3}}{3!} + \cdots$$

= $\sum_{k=0}^{\infty} \frac{z^{k}}{k!}$

The series always converges

The matrix exponential

Let A be a $(n \times n)$ matrix, by definition its exponential is a $(n \times n)$ matrix

$$e^{A} = I + A + \frac{A^{2}}{2!} + \frac{A^{3}}{3!} + \cdots$$

= $\sum_{k=0}^{\infty} \frac{A^{k}}{k!}$

The series always converges

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The state transition matrix (cont.)

The product of several matrices

The product of matrix A and B is only possible when the matrixes are compatible

• Number of columns of A must equal the number of rows of B

The same applies to the product of several matrixes

 $\underbrace{M}_{(m \times n)} = \underbrace{A_1}_{(m \times m_1)} \underbrace{A_2}_{(m_1 \times m_2)} \cdots \underbrace{A_{k-1}}_{(m_{k-2} \times m_{k-1})} \underbrace{A_k}_{(m_{k-1} \times n)}$

Powers of a matrix

Let A be an order-n square matrix, we want to define the k-th power of matrix A

The k-th power of matrix A is the n-order matrix A^k

$$A^k = \underbrace{A \times A \times \dots \times A}_{}$$

k times

Some special cases,

$$\stackrel{\rightsquigarrow}{\longrightarrow} A^{k=0} = I \\ \stackrel{\textstyle}{\longrightarrow} A^{k=1} = A$$

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The state transition matrix (cont.)

Definition

The state transition matrix

Consider a linear and time-invariant state-space model with $(N_x \times N_x)$ state matrix A

The state transition matrix of this system is given by the $(N_x \times N_x)$ matrix e^{At}

$$e^{At} = \underbrace{\frac{A^{0}t^{0}}{0!}}_{I} + \underbrace{\frac{A^{1}t^{1}}{1!}}_{At} + \underbrace{\frac{A^{2}t^{2}}{2!}}_{(A^{2}t^{2})/2!} + \cdots$$
$$= \sum_{k=0}^{\infty} \frac{A^{k}t^{k}}{k!}$$

The state transition matrix is well defined for any square matrix A

• (The series always converges)

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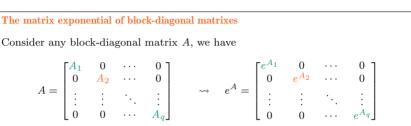
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The state transition matrix (cont.)

It is not convenient to determine the state transition matrix starting from its definition \rightsquigarrow One exception is when A is (block-)diagonal



The matrix exponential of diagonal matrixes (as special case)

For any diagonal $(n \times n)$ matrix A, we have

$$A = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & \lambda_n \end{bmatrix} \qquad \rightsquigarrow \quad e^A = \begin{bmatrix} e^{\lambda_1} & 0 & \cdots & 0 \\ 0 & e^{\lambda_2} & \cdots & 0 \\ \vdots & \vdots & \ddots & 0 \\ 0 & 0 & \cdots & e^{\lambda_n} \end{bmatrix}$$

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The state transition matrix (cont.)

Example

Consider a linear and time-invariant state-space model with (2×2) diagonal matrix A

$$A = \begin{bmatrix} -1 & 0\\ 0 & -2 \end{bmatrix}$$

We are interested in the corresponding state transition matrix

We have,

$$e^{At} = \begin{bmatrix} e^{(-1)t} & 0\\ 0 & e^{(-2)t} \end{bmatrix}$$

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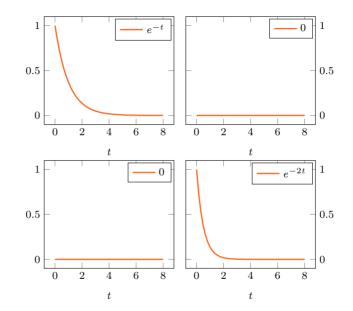
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Properties

We state without proof some fundamental results about a state transition matrix e^{At} \rightsquigarrow They are needed to derive Lagrange formula

Proposition

Derivative of the state transition matrix

Consider the state transition matrix e^{At} , we have,

$$\frac{\mathrm{d}}{\mathrm{d}t}e^{At} = Ae^{At}$$
$$= e^{At}A$$

By using the derivative property, we have that A commutes with e^{At} \rightsquigarrow (This result is important)

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Properties (cont.)

Proposition

Composition of two state transition matrices

Consider the two state transition matrices e^{At} and $e^{A\tau}$, we have

 $e^{At}e^{A\tau} = e^{A(t+\tau)}$

Proposition

Inverse of the state transition matrix

Let e^{At} be a state transition matrix, its inverse $(e^{At})^{-1}$ is matrix e^{-At}

$$e^{At}e^{-At} = e^{-At}e^{At}$$
$$= I$$

A state transition matrix e^{At} is always invertible (non-singular)

• Even if A were singular

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Properties (cont.)

Matrix inverse

Consider a square matrix A of order n

We define the **inverse** of A the square matrix of order n, A^{-1}

```
A^{-1}A = AA^{-1} = I
```

The inverse of matrix A exists if and only if A is non-singular

• When the inverse exists, it is also unique

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Sylvester's expansion

We determine the analytical expression of the state transition matrix e^{At}

- The procedure is known as Sylvester expansion
- (Does not require computing the infinite series)
- There are also other procedures (later)

Proposition

The Sylvester's expansion

Let A be a $(n \times n)$ matrix and let the corresponding state transition matrix be e^{At} . We have,

$$e^{At} = \beta_0(t)I + \beta_1(t)A + \beta_2(t)A^2 + \dots + \beta_{n-1}(t)A^{n-1}$$
$$= \sum_{i=0}^{n-1} \beta_i(t)A^i$$

The coefficients β_i of the expansion are appropriate functions of time

- \leadsto They can be determined by solving a set of linear equations
- \rightsquigarrow There is a finite number (n) of them

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Sylvester's expansion (cont.)

We show how to determine the coefficients when A has eigenvalues of multiplicity one

We will not consider the other cases, because rather involved and tedious to derive

- \rightsquigarrow Matrix A has complex eigenvalues (with multiplicity larger one)
- \rightsquigarrow Matrix A has complex eigenvalues (with multiplicity one)
- \rightsquigarrow Eigenvalues of A have multiplicity larger than one

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Sylvester's expansion (cont.)

Eigenvalues with multiplicity one

Let matrix A have distinct eigenvalues $\lambda_1, \lambda_2, \ldots, \lambda_n$

$$e^{At} = \sum_{i=0}^{n-1} \beta_i(t) A^i$$

= $\beta_0(t) I + \beta_1(t) A + \beta_2(t) A^2 + \dots + \beta_{n-1}(t) A^{n-1}$

The *n* unknown functions $\beta_i(t)$ are those that solve the system

$$\begin{cases} 1\beta_{0}(t) + \lambda_{1}\beta_{1}(t) + \lambda_{1}^{2}\beta_{2}(t) + \dots + \lambda_{1}^{n-1}\beta_{n-1}(t) = e^{\lambda_{1}t} \\ 1\beta_{0}(t) + \lambda_{2}\beta_{1}(t) + \lambda_{2}^{2}\beta_{2}(t) + \dots + \lambda_{2}^{n-1}\beta_{n-1}(t) = e^{\lambda_{2}t} \\ \dots \\ 1\beta_{0}(t) + \lambda_{n}\beta_{1}(t) + \lambda_{n}^{2}\beta_{2}(t) + \dots + \lambda_{n}^{n-1}\beta_{n-1}(t) = e^{\lambda_{n}t} \end{cases}$$

 $\sim \rightarrow$

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Sylvester's expansion (cont.)

Or, equivalently, in matrix form

 $V\beta = \eta$

• The vector of unknowns

 $\rightsquigarrow \quad \beta = \begin{bmatrix} \beta_0(t) & \beta_1(t) & \cdots & \beta_{n-1}(t) \end{bmatrix}^T$

• The coefficients matrix¹

• The known vector

 $^1\mathrm{A}$ matrix in this form is known to be a Vandermonde matrix.

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Sylvester's expansion (cont.)

$$\eta = \begin{bmatrix} e^{\lambda_1 t} & e^{\lambda_2 t} & \cdots & e^{\lambda_n t} \end{bmatrix}^T$$

The components of vector η are special functions of time, $e^{\lambda t}$ \rightsquigarrow Functions $e^{\lambda t}$ are the modes of matrix A

 \rightsquigarrow Mode $e^{\lambda t}$ associates with eigenvalue λ

Each element of e^{At} is a linear combination of such modes

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Sylvester expansion (cont.)

Example

Consider a (2×2) matrix A, we want to determine the state transition matrix e^{At}

$$A = \begin{bmatrix} -1 & 1\\ 0 & -2 \end{bmatrix}$$

Matrix A is triangular, the eigenvalues correspond to the diagonal elements Matrix A has 2 distinct eigenvalues $\rightsquigarrow \lambda_1 = -1$ $\rightsquigarrow \lambda_2 = -2$

To determine e^{At} , we write the system

$$\begin{cases} 1\beta_0(t) + \lambda_1\beta_1(t) = e^{\lambda_1 t} \\ 1\beta_0(t) + \lambda_2\beta_1(t) = e^{\lambda_2 t} \end{cases} \xrightarrow{\sim} \qquad \begin{cases} \beta_0(t) + (-1)\beta_1(t) = e^{(-1)\alpha_1 t} \\ \beta_0(t) + (-2)\beta_1(t) = e^{(-2)\alpha_1 t} \end{cases}$$

Sylvester's expansion (cont.)

By simple manipulation, we get

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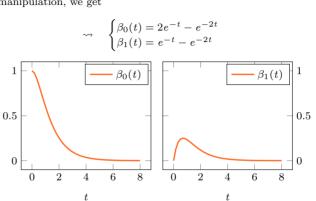
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Thus,

$$\begin{split} e^{At} &= \beta_0(t)I_2 + \beta_1(t)A = (2e^{-t} - e^{-2t}) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + (e^{-t} - e^{-2t}) \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-t} & (e^{-t} - e^{-2t}) \\ 0 & e^{-2t} \end{bmatrix} \end{split}$$

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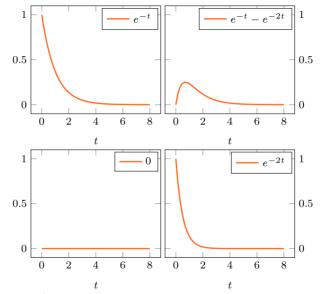
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Each element of e^{At} is a linear combination of the two system modes, e^{-t} and e^{-2t}

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Sylvester's expansion (cont.)

Eigenvalues and eigenvectors

Let $\lambda \in \mathcal{R}$ be some scalar and let $v \neq 0$ be some $(n \times 1)$ column vector

Consider a square matrix A of order n, suppose that the identify holds

```
Av = \lambda v
```

The scalar λ is called an **eigenvalue** of A

Vector v is the associated **eigenvector**

Consider a square matrix A of order n whose elements are real numbers

Matrix A has n (not necessarily distinct) eigenvalues $\lambda_1, \lambda_2, \ldots, \lambda_n$

- They can be real numbers or conjugate-complex pairs
- If $\lambda_i \neq \lambda_j$ for $i \neq j$, A has multiplicity one

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Sylvester's expansion (cont.)

Eigenvalues of triangular and diagonal matrices

Let matrix $A = \{a_{i,j}\}$ be a triangular or a diagonal matrix

• The eigenvalues of A are the n diagonal elements $\{a_{i,i}\}$

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Sylvester's expansion (cont.)

Characteristic polynomial

The characteristic polynomial of a square matrix A of order n

• The *n*-order polynomial in the variable s

 $P(s) = \det\left(sI - A\right)$

Computing eigenvalues and eigenvectors

The eigenvalues of matrix A of order n solve its characteristic polynomial \rightsquigarrow The roots of the equation $P(s) = \det (sI - A) = 0$

Let λ be an eigenvalue of matrix A

Each eigenvector v associated to it is a non-trivial solution to the system

 $(\lambda I - A)v = 0$

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Sylvester's expansion (cont.)

Systems of linear equations

Consider a system of *n* linear equations in *n* unknowns Ax = b $\rightsquigarrow A$ is a $(n \times n)$ matrix of coefficients

- $\rightarrow b$ is a $(n \times 1)$ vector of known terms
- $\rightsquigarrow x$ is a $(n \times 1)$ vector of unknowns

If A is non-singular, the system admits one and only one solution

If matrix A is singular, let M = [A|b] be a $[n \times (n+1)]$ matrix

- If rank(A) = rank(M), system has infinite solutions
- If rank(A) < rank(M), system has no solutions

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Sylvester's expansion (cont.)

Matrix rank

The rank of a $(m \times n)$ matrix A is equal to the number of columns (or rows) of the matrix that are linearly independent, rank(A)

Matrix kernel or null space

Consider a $(m \times n)$ matrix A, we define its null space or kernel

 $\ker(A) = \left\{ x \in \mathbb{R}^n | Ax = 0 \right\}$

It is the set of all vectors $x \in \mathbb{R}^n$ that left-multiplied by A produce the null vector

The set is a vector space, its dimension is called the **nullity** of matrix A, null(A)

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We can now prove the solution to the analysis problem for MIMO systems

• Lagrange formula

Theorem

Lagrange formula

Consider the state-space representation of a time-invariant linear system of order \boldsymbol{n}

The solution for $t \ge t_0$, for an initial state $x(t_0)$ and an input $u(t \ge t_0)$

$$\begin{aligned} x(t) &= e^{A(t-t_0)} x(t_0) + \int_{t_0}^t e^{A(t-\tau)} B u(\tau) \mathrm{d}\tau \\ y(t) &= C e^{A(t-t_0)} x(t_0) + C \int_{t_0}^t e^{A(t-\tau)} B u(\tau) \mathrm{d}\tau + D u(t) \end{aligned}$$

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Lagrange formula (cont.)

Proof

By left-multiplying the state equation $\dot{x}(t) = Ax(t) + Bu(t)$ by e^{-At} , we get

$$e^{-At}\dot{x}(t) = e^{-At}Ax(t) + e^{-At}Bu(t)$$

The resulting state equation can be rewritten,

$$e^{-At}\dot{x}(t) - e^{-At}Ax(t) = e^{-At}Bu(t)$$

Then, by using the result on the derivative of the state transition $matrix^2$,

$$\frac{\mathrm{d}}{\mathrm{d}t} \Big[e^{-At} x(t) \Big] = e^{-At} \dot{x}(t) - e^{-At} A x(t)$$
$$= e^{-At} B u(t)$$

²Derivative of the state transition matrix

$$\frac{\mathrm{d}}{\mathrm{d}t} \Big[e^{-At} x(t) \Big] = e^{-At} \Big[\frac{\mathrm{d}}{\mathrm{d}t} x(t) \Big] + \Big[\frac{\mathrm{d}}{\mathrm{d}t} e^{At} \Big] x(t) = e^{-At} \dot{x}(t) - e^{-At} A x(t).$$

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Lagrange formula (cont.)

$$\frac{\mathrm{d}}{\mathrm{d}t} \Big[e^{-At} x(t) \Big] = e^{-At} Bu(t)$$

By integrating between t_0 and t, we obtain

$$\left[e^{-A\tau}x(\tau)\right]_{t_0}^t = \int_{t_0}^t e^{-A\tau} Bu(\tau) \mathrm{d}\tau$$

That is,

$$e^{At}x(t) - e^{-At_0}x(t_0) = \int_{t_0}^t e^{-A\tau}Bu(t)d\tau$$

Thus,

$$e^{-At}x(t) = e^{-At_0}x(t_0) + \int_{t_0}^t e^{-A\tau}Bu(t)d\tau$$

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Lagrange formula (cont.)

$$e^{-At}x(t) = e^{-At_0}x(t_0) + \int_{t_0}^t e^{-A\tau}Bu(t)$$

The first Lagrange formula is obtained by multiplying both sides by e^{At}

$$\rightsquigarrow \quad x(t) = e^{A(t-t_0)}x(t_0) + \int_{t_0}^t e^{A(t-\tau)}Bu(\tau)\mathrm{d}\tau$$

The second formula is obtained by substituting x(t) in the output equation

$$y(t) = Cx(t) + Du(t)$$

$$\longrightarrow C\left[\underbrace{e^{A(t-t_0)}x(t_0) + \int_{t_0}^t e^{A(t-\tau)}Bu(\tau)d\tau}_{x(t)}\right] + Du(t)$$

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Force-free and forced evolution

$$x(t) = \underbrace{e^{A(t-t_0)}x(t_0)}_{x_u(t)} + \underbrace{\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)\mathrm{d}\tau}_{x_f(t)}$$

We can write the state solution (for $t \ge t_0$) as the sum of two terms

$$\rightsquigarrow$$
 $x(t) = x_u(t) + x_f(t)$

- \rightsquigarrow The force-free evolution of the state, $x_u(t)$
- \rightsquigarrow The forced evolution of the state, $x_f(t)$

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Force-free and forced evolution (cont.)

$$x(t) = \underbrace{e^{A(t-t_0)}x(t_0)}_{x_u(t)} + \underbrace{\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)d\tau}_{x_f(t)}$$

The force-free evolution of the state, from the initial condition $x(t_0)$

 $\rightsquigarrow \quad x_l(t) = e^{A(t-t_0)} x(t_0)$

 $\rightsquigarrow e^{A(t-t_0)}$ indicates the transition from $x(t_0)$ to x(t) \rightsquigarrow In the absence of contribution from the input

The forced evolution of the state

$$\rightsquigarrow \quad x_{f}(t) = \int_{t_{0}}^{t} e^{A(t-\tau)} Bu(\tau) \mathrm{d} au$$

- \rightsquigarrow The contribution of $u(\tau)$ to state x(t)
- \rightsquigarrow Through a weighting function, $e^{A(t-\tau)}B$

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Force-free and forced evolution (cont.)

$$y(t) = \underbrace{Ce^{A(t-t_0)}x(t_0)}_{y_u(t)} + \underbrace{C\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)\mathrm{d}\tau + Du(t)}_{y_f(t)}$$

We can write the output solution (for $t \ge t_0$) as the sum of two terms

$$\rightsquigarrow$$
 $y(t) = y_l(t) + y_f(t)$

- \rightsquigarrow The force-free evolution of the output, $y_u(t)$
- \rightsquigarrow The forced evolution of the output, $y_f(t)$

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$$y(t) = \underbrace{Ce^{A(t-t_0)}x(t_0)}_{y_u(t)} + \underbrace{C\int_{t_0}^t e^{A(t-\tau)}Bu(\tau)\mathrm{d}\tau + Du(t)}_{y_f(t)}$$

The force-free evolution of the output, from initial condition $y(t_0) = Cx(t_0)$

$$y_u(t) = Ce^{A(t-t_0)}x(t_0)$$
$$= Cx_u(t)$$

The **forced-evolution** of the output

$$\Rightarrow \quad y_f(t) = C \int_{t_0}^t e^{A(t-\tau)} Bu(\tau) d\tau + Du(t)$$
$$= C x_f(t) + Du(t)$$

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Free and forced evolution (cont.)

$$\underbrace{u(t)}_{y(t)} \xrightarrow{\dot{x}(t) = Ax(t) + Bu(t)}_{y(t) = Cx(t) + Du(t)} \underbrace{y(t)}_{y(t)}$$

System

Note that for $t_0 = 0$, we have

$$\begin{aligned} x(t) &= e^{At} x(0) + \int_0^t e^{A(t-\tau)} Bu(\tau) d\tau \\ y(t) &= C e^{At} x(0) + C \int_0^t e^{A(t-\tau)} Bu(\tau) d\tau + Du(t) d\tau \end{aligned}$$

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Example

Consider a linear time-invariant system with the state-space representation,

$$\begin{cases} \begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \end{bmatrix} &= \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u(t) \\ y(t) &= \begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

We want to determine the state and the output evolution for $t \ge 0$

- We consider the input signal $u(t) = 2\delta_{-1}(t)$
- We consider the initial state $x(0) = (3, 4)^T$

The state transition matrix for this state-space representation,

$$e^{At} = \begin{bmatrix} e^{-t} & (e^{-t} - e^{-2t}) \\ 0 & e^{-2t} \end{bmatrix}$$

We computed it earlier

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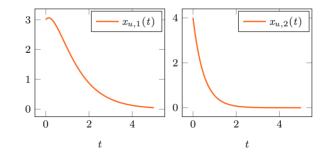
The force-free evolution of the state, for $t \ge 0$ $\Rightarrow r_{t}(t) = e^{At} r(0)$

$$x_u(t) = e^{At} x(0)$$

$$= \begin{bmatrix} e^{-t} & (e^{-t} - e^{-2t}) \\ 0 & e^{-2t} \end{bmatrix} \begin{bmatrix} 3 \\ 4 \end{bmatrix}$$

$$= \begin{bmatrix} (7e^{-t} - 4e^{-2t}) \\ 4e^{-2t} \end{bmatrix}$$

That is,



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Free and forced evolution (cont.)

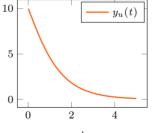
The force-free evolution of the output, for $t\geq 0$

$$y_u(t) = C x_u(t)$$

$$= \begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} (7e^{-t} - 4e^{-2t}) \\ 4e^{-2t} \end{bmatrix}$$

$$= 14e^{-t} - 4e^{-2t}$$

That is,



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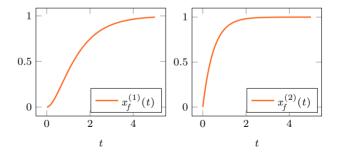
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Free and forced evolution (cont.)

The forced evolution of the state, for $t\geq 0$

 \rightarrow

$$\begin{aligned} x_f(t) &= \int_0^t e^{At} Bu(t-\tau) d\tau = \int_0^t \begin{bmatrix} e^{-\tau} & (e^{-\tau} - e^{-2\tau}) \\ 0 & e^{-2\tau} \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} 2d\tau \\ &= 2 \int_0^t \begin{bmatrix} (e^{-\tau} - e^{-2\tau}) \\ e^{-2\tau} \end{bmatrix} d\tau = 2 \begin{bmatrix} \int_0^t (e^{-\tau} - e^{-2\tau}) d\tau \\ \int_0^t e^{-2t} d\tau \end{bmatrix} \\ &= 2 \begin{bmatrix} (1 - e^{-t}) - 1/2(1 - e^{-2t}) \\ 1/2(1 - e^{-2t}) \end{bmatrix} = \begin{bmatrix} (1 - 2e^{-t} + e^{-2t}) \\ (1 - e^{-2t}) \end{bmatrix} \end{aligned}$$



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Free and forced evolution (cont.)

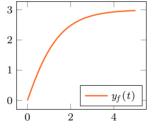
 \rightarrow

Since D = 0, the forced evolution of the output for $t \ge 0$

$$y_f(t) = Cx_f(t)$$

= $\begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} (1 - 2e^{-t} + e^{-2t}) \\ (1 - e^{-2t}) \end{bmatrix}$
= $3 - 4e^{-t} + e^{-2t}$

That is,



t

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The form of the state space representation depends on the choice of state variables

• The choice is not unique, even when we are coming from a physical model

There is an infinite number of different representations of the same system

- They are all related by a similarity transformation
- These transformations allow flexibility in the analysis
- We can change to easier system representations

The state matrix can be set to a canonical form

- → Diagonal form
- → Jordan form

 $\rightsquigarrow \cdots$

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Similarity transformation (cont.)

Definition

Similarity transformation

Consider the state-space representation of a linear time-invariant system of order N_x

$$\begin{array}{c} u(t) \\ \hline \dot{x}(t) = Ax(t) + Bu(t) \\ y(t) = Cx(t) + Du(t) \end{array} \begin{array}{c} y(t) \\ \hline \dot{x}(t) = Ax(t) + Bu(t) \\ y(t) = Cx(t) + Du(t) \end{array}$$

System

- x(t) and $\dot{x}(t)$, state vector and its derivative (N_x components)
- u(t), input vector (N_u components)
- y(t), output vector (N_y components)

Let vector z(t) be related to x(t) by some linear transformation P, x(t) = Pz(t)

P is any $(N_x \times N_x)$ non-singular matrix of constants (its inverse always exists)

• Because of non-signularity, we have $z(t) = P^{-1}x(t)$

The transformation/matrix P is called a similarity transformation/matrix

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Similarity transformation (cont.)

Proposition

Similar representations

Consider the state-space representation of a linear time-invariant system of order N_x

$$\begin{array}{c} u(t) \\ \hline \dot{x}(t) = Ax(t) + Bu(t) \\ y(t) = Cx(t) + Du(t) \end{array} \begin{array}{c} y(t) \\ \hline \\ y(t) = Cx(t) + Du(t) \end{array} \end{array} \qquad \begin{cases} \dot{x}(t) = Ax(t) + Bu(t) \\ y(t) = Cx(t) + Du(t) \\ \hline \\ y(t) = Cx(t) + Du(t) \end{array}$$

Let P be some similarity transformation matrix such that x(t) = Pz(t)Vector $z(t) = P^{-1}x(t)$ satisfies the new state-space representation

$$u(t) \quad \dot{x}(t) = A'x(t) + B'u(t) \quad y(t)$$

$$y(t) = C'x(t) + D'u(t) \quad (x + D'u(t))$$

$$y(t) = C'x(t) + D'u(t) \quad (x + D'u(t))$$

$$y(t) = C'z(t) + D'u(t)$$

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Similarity transformation (cont.)

Proof

By taking the time-derivative of the state vector x(t) = Pz(t), we have

$$\rightsquigarrow \dot{x}(t) = P\dot{z}(t)$$

By substituting x(t) and $\dot{x}(t)$ into the state-space representation,

$$\stackrel{\sim}{\rightarrow} \quad \begin{cases} \underbrace{P\dot{z}(t)}_{\dot{x}(t)} = A \underbrace{Pz(t)}_{x(t)} + Bu(t) \\ \underbrace{y(t)}_{y(t)} = C \underbrace{Pz(t)}_{x(t)} + Du(t) \\ \underbrace{y(t)}_{x(t)} = C \underbrace{Pz(t)}_{x(t)} + Du(t) \end{cases}$$

Pre-multiply the state equation by P^{-1} , to complete the proof

$$P^{-1}P\dot{z}(t) = P^{-1}APz(t) + P^{-1}Bu(t)$$
$$P^{-1}y(t) = P^{-1}CPz(t) + P^{-1}Du(t)$$

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$$P^{-1}P\dot{z}(t) = P^{-1}APz(t) + P^{-1}Bu(t)$$
$$P^{-1}y(t) = P^{-1}CPz(t) + P^{-1}Du(t)$$

For the state equation, we have

$$\underbrace{P^{-1}P}_{I}\dot{z}(t) = \underbrace{P^{-1}AP}_{A'}z(t) + \underbrace{P^{-1}B}_{B'}u(t)$$

For the measurements, we have

$$\underbrace{\frac{PP^{-1}}{I}y(t) = \underbrace{PP^{-1}}_{I}CPz(t) + \underbrace{PP^{-1}}_{I}Du(t)}_{CP}z(t) + \underbrace{D}_{D'}u(t)$$

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U

$$\underbrace{\dot{x}(t) = A'x(t) + B'u(t)}_{y(t) = C'x(t) + D'u(t)} \underbrace{y(t)}_{y(t) = C'z(t) + D'u(t)} \left\{ \begin{aligned} \dot{z}(t) = A'z(t) + B'u(t) \\ y(t) = C'z(t) + D'u(t) \end{aligned} \right\}$$

We obtained a different state-space representation of the same dynamical system

- Input u(t) and output y(t) are left unchanged (problem data)
- We defined a new (transformed) state variables, z(t)

There is an infinite number of non-singular matrixes P that could be used \rightsquigarrow Thus, there is also an infinite number of equivalent representations $\rightsquigarrow A' = P^{-1}AP$ $\rightsquigarrow B' = P^{-1}B$ $\rightsquigarrow C' = CP$ $\rightsquigarrow D' = D$

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Similarity transformation (cont.)

Example

Consider a linear time-invariante system with state-space representation $\{A,B,C,D\}$

$$\begin{cases} \begin{bmatrix} \dot{x_1}(t) \\ \dot{x_2}(t) \end{bmatrix} = \overbrace{\begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix}}^{A} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \overbrace{\begin{bmatrix} 0 \\ 1 \end{bmatrix}}^{B} u(t) \\ \begin{bmatrix} y_1(t) \\ y_2(t) \end{bmatrix} = \underbrace{\begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix}}_{C} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \underbrace{\begin{bmatrix} 1.5 \\ 0 \end{bmatrix}}_{D} u(t) \end{cases}$$

Consider the similarity transformation of the state using some matrix ${\cal P}$

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \underbrace{\begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}}_{P} \begin{bmatrix} z_1(t) \\ z_2(t) \end{bmatrix}$$

What is the $\{A', B', C', D'\}$ state-space representation for state z(t)

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Similarity transformation (cont.)

We are given the similarity transformation P,

$$P = \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}$$

We compute its inverse,

$$P^{-1} = \begin{bmatrix} 0 & 1\\ 1 & -1 \end{bmatrix}$$

Since
$$z(t) = P^{-1}x(t)$$
, we have

$$\begin{bmatrix} z_1(t) \\ z_2(t) \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$
$$= \begin{bmatrix} x_2(t) \\ x_1(t) - x_2(t) \end{bmatrix}$$

The second component of z is the difference between first and second component of x \rightsquigarrow The first component of z simply equals the second component of x

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Similarity transformation (cont.)

We conclude by calculating the resulting state-space representation

$$\begin{aligned} A' &= P^{-1}AP \\ &= \begin{bmatrix} 0 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 0 & -1 \\ -2 & 0 \end{bmatrix} = \begin{bmatrix} -2 & 0 \\ 2 & -1 \end{bmatrix} \\ B' &= P^{-1}B \\ &= \begin{bmatrix} 0 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix} \\ C' &= CP \\ &= \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 3 & 2 \\ 2 & 0 \end{bmatrix} \\ D' &= D \\ &= \begin{bmatrix} 1.5 \\ 0 \end{bmatrix} \end{aligned}$$

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Similarity transformation (cont.)

Proposition

Similarity and state transition matrix

Consider the state matrix $A' = P^{-1}AP$ from some similarity transformation P

The corresponding state transition matrix,

$$e^{A't} = P^{-1}e^{At}P$$

Proof

Note that

$$(A')^{k} = \underbrace{(P^{-1}AP) \cdot (P^{-1}AP) \cdots (P^{-1}AP)}_{k \text{ times}} = P^{-1}\underbrace{AA \cdots A}_{k \text{ times}} P = P^{-1}A^{k}P$$

Thus, by definition

$$e^{A't} = \sum_{k=0}^{\infty} \frac{(A')^k t^k}{k!} = \sum_{k=0}^{\infty} \frac{(P^{-1}A^k P)t^k}{k!} = P^{-1} \Big(\sum_{k=0}^{\infty} \frac{A^k t^k}{k!}\Big) P = P^{-1}e^{At}P$$

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Similarity transformation (cont.)

We show how two similar state-space representations describe the same IO relation

Proposition

Invariance of the IO relationship under similarity

Consider two similar state-space representations of a linear time-invariant system $% \left({{{\mathbf{x}}_{i}}} \right)$

- $\rightsquigarrow \{A, B, C, D\} \text{ and } \{A', B', C', D'\}$
- $\rightsquigarrow P$ is the transformation matrix

Suppose that the system be subjected to some known input $\rightsquigarrow u(t)$

The two representations produce the same forced response $\rightsquigarrow \ y_f(t)$

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Similarity transformation (cont.)

${\bf Proof}$

Consider the Lagrange formula

The forced response of the second representation due to input u(t)

$$y_{f}(t) = C' \int_{t_{0}}^{t} e^{A'(t-\tau)} B'u(\tau) d\tau + Du(t)$$

= $CP \int_{t_{0}}^{t} \underbrace{P^{-1}e^{A(t-\tau)}P}_{e^{A'(t-\tau)}} \underbrace{P^{-1}B}_{B'} u(\tau) d\tau + Du(t)$
= $C \int_{t_{0}}^{t} e^{A(t-\tau)} Bu(\tau) d\tau + Du(t)$

This response corresponds to the one of the original representation

$$y_f(t) = C \int_{t_0}^t e^{A(t-\tau)} Bu(\tau) \mathrm{d}\tau + Du(t)$$

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Similarity transformation (cont.)

Proposition

Invariance of the eigenvalues under similarity transformations

Matrix A and $P^{-1}AP$ have the same characteristic polynomial

Proof

The characteristic polynomial of matrix A^\prime

$$\det (\lambda I - A') = \det (\lambda I - P^{-1}AP)$$
$$= \det (\lambda \underbrace{P^{-1}P}_{I} - P^{-1}AP)$$
$$= \det [P^{-1}(\lambda I - A)P]$$
$$= \det (P^{-1}) \det (\lambda I - A) \det (P)$$
$$= \det (\lambda I - A)$$

The last equality is obtained from $det(P^{-1})det(P) = 1$

A and A^\prime share the same characteristic polynomial

 $\rightsquigarrow\,$ Thus, also the eigenvalues are the same

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Two similar representations have the same modes, the modes characterise the dynamics

The modes are therefore independent of the representation

Similarity transformation (cont.)

 \rightsquigarrow This is important

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Similarity transformation (cont.)

lxample

Consider two similar state-space representations of a linear time-invariant system

$$A = \begin{bmatrix} -1 & 1\\ 0 & -2 \end{bmatrix}$$
$$A' = \begin{bmatrix} -2 & 0\\ 2 & -1 \end{bmatrix}$$

The similarity transformation matrix

$$P = \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}$$

We are interested in the eigenvalues and modes of the system

Matrix A and A have two eigenvectors

- $\lambda_1 = -1$
- $\lambda_2 = -2$

The system modes are e^{-t} and e^{-2t}

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Diagonalisation

We consider a special similarity transformation P, we seek for a diagonal matrix $A' \rightarrow A$ state-space representation with a diagonal state matrix

 \rightsquigarrow Diagonal canonical form

 $\rightsquigarrow \Lambda = A' = P^{-1}AP$

Consider the linear time-invariant system with a single input (and, say, single output)

$$\begin{bmatrix} \dot{x}_1(t)\\ \dot{x}_2(t)\\ \vdots\\ \dot{x}_n(t) \end{bmatrix} = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0\\ 0 & \lambda_2 & \cdots & 0\\ \vdots & \vdots & \ddots & \vdots\\ 0 & 0 & \cdots & \lambda_n \end{bmatrix} \begin{bmatrix} x_1(t)\\ x_2(t)\\ \vdots\\ x_n(t) \end{bmatrix} + \begin{bmatrix} b_1\\ b_2\\ \vdots\\ b_n \end{bmatrix} u(t)$$

The evolution of the i-th component of the state vector

$$\rightsquigarrow \quad \dot{x}_i(t) = \lambda_i x_i(t) + b_i u(t)$$

State derivatives are not related to other components

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Diagonalisation (cont.)

We can understand a system with diagonal matrix ${\cal A}$ as a collection of sub-systems

- $\rightsquigarrow\,$ Each sub-system is described by a single state component
- $\rightsquigarrow\,$ Each state component evolves independently
- \rightsquigarrow The representation is **decoupled**
- $\rightsquigarrow N_x$ first-order subsystems

The characteristic polynomial of the system for the i-th component

$$\rightsquigarrow$$
 $P_i(s) = (s - \lambda_i)$

This subsystem has mode $e^{-\lambda_i t}$

We show how to determine a similarity transformation that leads to a diagonal form

• This can be understood as a somehow special similarity transformation

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Diagonalisation (cont.)

Definition

Modal matrix

Consider a system in state-space representation with $(N_x \times N_x)$ matrix A

- Let v_1, v_2, \ldots, v_n be a set of all the eigenvectors of matrix A
- Suppose that they correspond to eigenvalues $\lambda_1, \lambda_2, \ldots, \lambda_n$

Suppose that eigenvectors in this set are also linearly independent

We define the **modal matrix** of A as the $(N_x \times N_x)$ matrix V

 $V = \begin{bmatrix} v_1 | v_2 | \cdots | v_n \end{bmatrix}$

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Diagonalisation (cont.)

If a matrix A has N_x distinct eigenvalues λ, then the modal matrix A always exists
Its N_x eigenvectors v are linearly independent

Distinct eigenvalues

Let A be a n-order matrix whose n eigenvalues $\lambda_1, \lambda_2, \ldots, \lambda_n$ are distinct

Then, there is a set of n linearly independent eigenvectors

• Vectors v_1, v_2, \ldots, v_n form a basis for \mathcal{R}^n

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Diagonalisation (cont.)

Example

Consider a state-space representation of a linear time-invariant system with matrix A

$$A = \begin{bmatrix} 2 & 1 \\ 3 & 4 \end{bmatrix}$$

We are interested in the modal matrix $\,V$ of A

The eigenvalues and eigenvectors of A

$$\stackrel{\rightsquigarrow}{\longrightarrow} \lambda_1 = 1 \text{ and } v_1 = \begin{bmatrix} 1 & -1 \end{bmatrix}^T$$
$$\stackrel{\rightsquigarrow}{\longrightarrow} \lambda_2 = 5 \text{ and } v_2 = \begin{bmatrix} 1 & 3 \end{bmatrix}^T$$

The modal matrix V,

$$V = \begin{bmatrix} v_1 | v_2 \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ -1 & 3 \end{bmatrix}$$

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$$V = \begin{bmatrix} v_1 | v_2 \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ -1 & 3 \end{bmatrix}$$

It is important to note that the eigenvectors are determined up to a scaling constant

- (Plus, the ordering of the eigenvalues is arbitrary)
- There can be more than one modal matrix

These two modal matrices of matrix A are equivalent

$$V' = \begin{bmatrix} 2v_1 | 3v_2 \end{bmatrix} = \begin{bmatrix} 2 & 3\\ -2 & 9 \end{bmatrix}$$
$$V'' = \begin{bmatrix} v_2 | v_1 \end{bmatrix} = \begin{bmatrix} 1 & 1\\ 3 & -1 \end{bmatrix}$$

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Consider a matrix A with some eigenvalues λ that have multiplicity ν larger than one

• The modal matrix V exists if and only if to each eigenvalue λ_i with multiplicity ν_i is possible to associate ν_i linearly independent eigenvectors $\{v_{i,1}, v_{i,2}, \ldots, v_{i,\nu_i}\}$

This is not always possible

But, ...

If a matrix admits a modal matrix, then it can be diagonalised

• (This is what matters to us)

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Example

Consider a state space representation of a linear time-invariant system with matrix A

$$A = \begin{bmatrix} 2 & 0\\ 0 & 2 \end{bmatrix}$$

Its only eigenvalue $\lambda=2$ has multiplicity $\nu=2$

Its eigenvectors are obtained by solving the system $[\lambda I - A]v = 0$

$$\begin{bmatrix} 2I - A \end{bmatrix} v = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad \rightsquigarrow \quad \begin{cases} 0 = 0 \\ 0 = 0 \end{cases}$$

We can choose any two linearly independent eigenvectors for λ

• As the equation is satisfied for any value of a and b

A modal matrix with the eigenvectors from the canonical basis

$$\rightsquigarrow \quad V = \begin{bmatrix} v_1 | v_2 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

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Consider a state space representation of a linear time-invariant system with matrix A

$$A = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix}$$

Its only eigenvalue $\lambda = 2$ has multiplicity $\nu = 2$

Its eigenvectors are obtained by solving the system $[\lambda I - A]v = 0$

$$[2I - A]v = \begin{bmatrix} 0 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad \rightsquigarrow \quad \begin{cases} -b = 0 \\ 0 = 0 \end{cases}$$

As b = 0, we can choose only one linearly independent eigenvector for λ

$$v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Matrix A does not admit a modal matrix

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Proposition

Diagonalisation

Consider a state-space representation of a linear time-invariant system with matrix ALet $\lambda_1, \lambda_2, \ldots, \lambda_n$ be its eigenvalues and $V = [v_1|v_2|\cdots|v_n]$ one of its modal matrices

Let Λ be the state matrix transformed according to $\Lambda = V^{-1}AV$ \rightsquigarrow Λ is diagonal

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Example

Consider a linear time-invariant system with matrixes $\{A,B,C,D\}$

$$\begin{cases} \begin{bmatrix} \dot{x_1}(t) \\ \dot{x_2}(t) \end{bmatrix} = \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u(t) \\ \begin{bmatrix} y_1(t) \\ y_2(t) \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 1.5 \\ 0 \end{bmatrix} u(t)$$

We are interested in a diagonal representation by similarity

We can compute the eigenvalues and eigenvectors of A

• $\lambda_1 = -1$ and $v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$ • $\lambda_2 = -2$ and $v_2 = \begin{bmatrix} 1 \\ -1 \end{bmatrix}$

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Then, we can determine a modal matrix and its inverse

$$V = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix}$$
$$V^{-1} = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix}$$

From the similarity transformation expressions, we get

$$\begin{aligned} A' &= V^{-1}AV = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix} \\ &= \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} -1 & -2 \\ 0 & 2 \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} = \Lambda \\ B' &= V^{-1}B = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ -1 \end{bmatrix} \\ C' &= CV = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 0 & -2 \end{bmatrix} \\ D' &= D = \begin{bmatrix} 1.5 \\ 0 \end{bmatrix} \end{aligned}$$

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State transition matrix by diagonalisation

We show a procedure alternative to Sylvester's formula for the state transition matrix

- We assume a linear time-invariant state-space system representation
- We assume that the state matrix A can be diagonalised

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State transition matrix by diagonalisation

Consider a $(n \times n)$ state matrix A and let $\lambda_1, \lambda_2, \ldots, \lambda_n$ be its eigenvalues

Suppose that A admits the modal matrix V

We have for the state transition matrix

$$e^{At} = V e^{\Lambda t} V^{-1} = V \begin{bmatrix} e^{\lambda_1 t} & 0 & \cdots & 0 \\ 0 & e^{\lambda_2 t} & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & e^{\lambda_n t} \end{bmatrix} V^{-1}$$

Because we have a diagonal state matrix

$$\Lambda = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_n \end{bmatrix}$$

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Proof

We have shown that the identity holds (see similarity and state transition matrices³) $e^{\Lambda t} = V^{-1} e^{At} V$

To complete the proof, multiply both sides by V on the left and by V^{-1} on the right

³Given $A' = P^{-1}AP$, we have $e^{A't} = P^{-1}e^{At}P$.

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Example

Consider a linear time-invariant system with matrixes $\{A, B, C, D\}$

$$\begin{cases} \begin{bmatrix} \dot{x_1}(t) \\ \dot{x_2}(t) \end{bmatrix} = \begin{bmatrix} -1 & 1 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u(t) \\ \begin{bmatrix} y_1(t) \\ y_2(t) \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 1.5 \\ 0 \end{bmatrix} u(t) \end{cases}$$

We are interested in computing the state transition matrix e^{At}

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State transition matrix by diagonalisation (cont.)

We have already computed the modal matrix of A and its inverse, V and V^{-1}

$$V = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix}$$
$$V^{-1} = \begin{bmatrix} 1 & 1 \\ 0 & -1 \end{bmatrix}$$

Thus, we have

$$e^{At} = V \begin{bmatrix} e^{\lambda_1 t} & 0\\ 0 & e^{\lambda_2 t} \end{bmatrix} V^{-1} = \begin{bmatrix} 1 & 1\\ 0 & -1 \end{bmatrix} \begin{bmatrix} e^{-t} & 0\\ 0 & e^{-2t} \end{bmatrix} \begin{bmatrix} 1 & 1\\ 0 & -1 \end{bmatrix}$$
$$= \begin{bmatrix} 1 & 1\\ 0 & -1 \end{bmatrix} \begin{bmatrix} e^{-t} & e^{-t}\\ 0 & -e^{-2t} \end{bmatrix} = \begin{bmatrix} e^{-t} & (e^{-t} - e^{-2t})\\ 0 & e^{-2t} \end{bmatrix}$$

This is the same result we determined by using the Sylvester expansion