

CHEM-LV03
2022

Informal
definition

Examples

Fokker-Planck

Definition

Derivation



Aalto University

Stochastic differential equations

Stochastic analysis and simulation of reactive and diffusive systems

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Informal definition

SDEs | Informal definition

Consider the evolution of some variable $x(t) \in \mathbb{R}$ according to a differential equation

$$\frac{dx(t)}{dt} = f(x(t), t), \quad \text{with } f : \mathbb{R} \times [0, \infty) \rightarrow \mathbb{R}$$

Given the initial condition $x(t=0) = x_0$, we solve the ordinary differential equation

- The solution $(x(t))_{t \geq 0}$ exists if f is a ‘well-behaved’ function
- (Conditions for existence and uniqueness can be stated)

We can formally re-write the ordinary differential equation

$$dx(t) = f(x(t), t) dt$$

The infinitesimal change of variable x during $[t, t + dt)$

$$\begin{aligned} x(t + dt) - x(t) &= dx(t) \\ &= f(x(t), t) dt \end{aligned}$$

To solve the equation, we have the simple recursion

$$x(t + dt) = x(t) + f(x(t), t) dt$$

SDEs | Informal definition (cont.)

$$x(t + dt) = x(t) + f(x(t), t) dt$$

To compute the solution of the ODE, we consider a small time interval of duration Δt

$$x(t + \Delta t) = x(t) + f(x(t), t) \Delta t$$

Then, given some initial condition $x(0) = x_0$, we have the collection $\{x(k\Delta t)\}_{k=0}^K$

$$x(0 + 1\Delta t) = x(0\Delta t) + f(x(0\Delta t), 0\Delta t) \Delta t$$

$$x(0 + 2\Delta t) = x(1\Delta t) + f(x(1\Delta t), 1\Delta t) \Delta t$$

$$\dots = \dots$$

$$x(0 + k\Delta t) = x(k\Delta t) + f(x(k\Delta t), k\Delta t) \Delta t$$

$$\dots = \dots$$

The iterative scheme is the (explicit) Euler’s method to approximate ODE’s solutions

- The approximation gets better, the smaller the duration of the interval Δt

SDEs | Informal definition (cont.)

We can use an informal definition of a stochastic differential equation (SDE) as an ordinary differential equation with an additional term describing stochastic fluctuations

$$X(t + \Delta t) = X(t) + f(X(t), t) \Delta t + \underbrace{g(X(t), t) \sqrt{\Delta t} \mathcal{N}(0, 1)}_{\text{additional term}}$$

We have the deterministic functions

$$\begin{cases} f : \mathbb{R} \times [0, 8) \rightarrow \mathbb{R} \\ g : \mathbb{R} \times [0, 8) \rightarrow \mathbb{R} \end{cases}$$

Function $g(\cdot)$ characterises the strength of the additive stochastic term, scaled by $\sqrt{\Delta t}$

Normally distributed numbers with zero-mean and unit-variance are easily simulated

$$\xi(t) \sim \mathcal{N}(0, 1)$$

In terms of the Wiener increment $\Delta W(t)$ of the standard Brownian motion $W(t)$,

$$X(t + \Delta t) = X(t) + f(X(t), t) \Delta t + \underbrace{g(X(t), t) \sqrt{\Delta t} \xi(t)}_{\Delta W(t)}$$

SDEs | Informal definition (cont.)

Given an initial condition $x(0) \sim p(X(0))$, a $\xi(0) \sim \mathcal{N}(0, 1)$, and small interval Δt

$$x(0 + 1\Delta t) = x(0\Delta t) + f(x(0\Delta t), 0\Delta t) \Delta t + \underbrace{g(x(0\Delta t), 0\Delta t) \sqrt{\Delta t} \xi(0\Delta t)}_{\Delta W(0\Delta t)}$$

$$x(0 + 2\Delta t) = x(1\Delta t) + f(x(1\Delta t), 1\Delta t) \Delta t + \underbrace{g(x(1\Delta t), 1\Delta t) \sqrt{\Delta t} \xi(1\Delta t)}_{\Delta W(1\Delta t)}$$

... = ...

$$x(0 + k\Delta t) = x(k\Delta t) + f(x(k\Delta t), k\Delta t) \Delta t + \underbrace{g(x(k\Delta t), k\Delta t) \sqrt{\Delta t} \xi(k\Delta t)}_{\Delta W(k\Delta t)}$$

... = ...

As a result of a single realisation of the recursion, we get the collection $\{x(k\Delta t)\}_{k=1}^K$

- The realisation scheme is the Euler-Maruyama method for simulating SDEs

SDEs | Informal definition (cont.)

$$X(t + dt) = X(t) + f(X(t), t) dt + g(X(t), t) \underbrace{\sqrt{dt}\xi(t)}_{dW(t)}$$

That is,

$$dX(t) = f(X(t), t) dt + g(X(t), t) dW(t)$$

The solution of the differential equation is a stochastic process $(X(t))_{t \geq 0}$

$$X(t) = X(0) + \int_0^t f(x(s), s) ds + \int_0^t g(x(s), s) dW_s$$

Both function $f(\cdot, \cdot)$ and $g(\cdot, \cdot)$ are continuous functions of x and t

SDEs | Informal definition (cont.)

$$X(t) = X(t_0) + \int_0^t f(x(s), s) ds + \int_0^t g(x(s), s) dW_s$$

We can re-write the solution of the stochastic differential equation using summations,

$$X(t) = X(0) + \lim_{\Delta t_k \rightarrow 0} \sum_{0 < t_k < t} f(X(t'_k), t'_k) \Delta t_k + \lim_{\Delta s_k \rightarrow 0} \sum_{0 < s_k < t} g(X(s'_k), s'_k) \Delta W_{s_k}$$

We introduced two partitions $\{t_k\}_{k=0}^K$ and $\{s_k\}_{k=0}^K$ of the $[0, t]$ interval,

$$\begin{array}{cccccccccccc} \underbrace{t_0}_{0} & \cdots & t_1 & \cdots & \cdots & \cdots & \underbrace{t_{k-1} \cdots t_k}_{\Delta t_k} & \cdots & \underbrace{t_{k+1} \cdots t_{k+1}}_{\Delta t_k} & \cdots & \cdots & t_{K-1} & \cdots & \underbrace{t_K}_t \\ \underbrace{s_0}_{0} & \cdots & s_1 & \cdots & \cdots & \cdots & \underbrace{s_{k-1} \cdots s_k}_{\Delta s_k} & \cdots & \underbrace{s_{k+1} \cdots s_{k+1}}_{\Delta s_k} & \cdots & \cdots & s_{K-1} & \cdots & \underbrace{s_K}_t \end{array}$$

We introduced two interval points $t'_k \in \Delta t_k$ and $s'_k \in \Delta s_k$,

$$\begin{aligned} t_k &\leq t'_k \leq t_{k+1} \\ s_k &\leq s'_k \leq s_{k+1} \end{aligned}$$

We also used $\Delta W_{s_k} = W(s_{k+1}) - W(s_k)$

SDEs | Informal definition (cont.)

$$X(t) = X(0) + \lim_{\Delta t_k \rightarrow 0} \sum_{0 < t_k < t} f(X(t'_k), t'_k) \Delta t_k + \underbrace{\lim_{\Delta s_k \rightarrow 0} \sum_{0 < s_k < t} g(X(s'_k), s'_k) \Delta W_{s_k}}_{\text{A stochastic process } I(g(\cdot, \cdot), t)}$$

For $\Delta t_k \rightarrow 0$, the first sum converges to the Riemann integral regardless of $t'_k \in \Delta t_k$

$$\rightsquigarrow \int_0^t f(X(s), s) ds$$

For $\Delta s_k \rightarrow 0$, the value of the second sum may depend on the choice of $s'_k \in \Delta s_k$

SDEs | Informal definition (cont.)

We can verify the dependence of the second integral on the actual choice of $s'_k \in \Delta s_k$

Consider the following integral

$$\int_0^t W(s) dW_s \simeq \lim_{\Delta s_k \rightarrow 0} \sum_{0 < s_k < t} W(s'_k) \Delta W_{s_k}$$

Introduce a partition $\{s_k\}_{k=0}^K$ of $[0, t]$,

$$\underbrace{s_0}_{0} \cdots s_1 \cdots \cdots \underbrace{s_{k-1} \cdots s_k}_{\Delta s_k} \cdots \underbrace{s_{k+1} \cdots s_{k+1}}_{\Delta s_k} \cdots \cdots s_{K-1} \cdots \underbrace{s_K}_t$$

Introduce the points $s'_k \in \Delta s_k$,

$$s_k \leq s'_k \leq s_{k+1}$$

Again, we set

$$\Delta W_{s_k} = W(s_{k+1}) - W(s_k)$$

We will consider three separate cases

- $s'_k = s_k$
- s'_k such that $W_{s'_k} = (W_{s_{k+1}} + W_{s_k})/2$

SDEs | Informal definition (cont.)

$$\int_0^t W(s) dW_s \simeq \lim_{\Delta s_k \rightarrow 0} \sum_{0 < s_k < t} W(s'_k) \Delta W_{s_k}$$

Case I: $s'_k = s_k$

$$\begin{aligned} \left\langle \int_0^t W_s dW_s \right\rangle &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{0 < s'_k < t} W(s'_k) \Delta W_{s_k} \right\rangle \\ &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s'_k \in \{s_k\}_{k=0}^{K-1}} W(s'_k) \Delta W_{s_k} \right\rangle \\ &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} W(s_k) \Delta W_{s_k} \right\rangle \\ &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} W(s_k) (W(s_k+1) - W(s_k)) \right\rangle \\ &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} \left[\frac{1}{2} (W(s_k+1) + W(s_k)) - \frac{1}{2} (W(s_k+1) - W(s_k)) \right] (W(s_k+1) - W(s_k)) \right\rangle \\ &= \frac{1}{2} \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} [W^2(s_k+1) - W^2(s_k)] \right\rangle - \frac{1}{2} \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} [W(s_k+1) - W(s_k)]^2 \right\rangle \\ &= \frac{1}{2} (\langle W^2(t) \rangle - \langle W^2(0) \rangle) - \frac{1}{2} t \\ &= 0 \end{aligned}$$

SDEs | Informal definition (cont.)

$$\int_0^t W(s) dW_s \simeq \lim_{\Delta s_k \rightarrow 0} \sum_{0 < s_k < t} W(s'_k) \Delta W_{s_k}$$

Case III: s'_k such that $W_{s'_k} = (W_{s_k+1} + W_{s_k})/2$

$$\begin{aligned} \left\langle \int_0^t W_s dW_s \right\rangle &= \left\langle \lim_{\Delta s_k \rightarrow 0} \sum_{s_k=0}^{K-1} (W(s_k+1) - W(s_k)) (W(s_k+1) - W(s_k)) \right\rangle \\ &= \frac{1}{2} \left\langle \sum_{\Delta s \rightarrow 0} \sum_{s_k=0}^{K-1} (W(s_k+1)^2 - W(s_k)^2) \right\rangle \\ &= \frac{1}{2} (\langle W(t)^2 \rangle - \langle W(0)^2 \rangle) \\ &= \frac{t}{2} \end{aligned}$$

SDEs | Informal definition (cont.)

We show that there are different definitions of the stochastic integral

$$I(g(\cdot, \cdot), t) = \int_0^t g(X(s), s) dW_s$$

There are two main types of definition of the stochastic integral

- The Itô integral

$$\rightsquigarrow s'_k = s_k$$

- The Stratonovich integral

$$\rightsquigarrow s'_k \text{ such that } g(X(s'_k), s'_k) = 1/2 [g(X(s_k), s_k) + g(X(s_k + 1), s_k + 1)]$$

SDEs | Informal definition (cont.)

Accordingly, we also have two interpretations of the stochastic differential equation

- The Itô equation

$$dX(t) = f(X(t), t) dt + g(X(t), t) dW(t)$$

- The Stratonovich equation

$$dX(t) = f(X(t), t) dt + g(X(t), t) \circ dW(t)$$

The Stratonovich equation is equivalent to the Itô equation of the form

$$dX(t) = \left[f(X(t), t) + \frac{1}{2} \frac{\partial g(X(t), t)}{\partial x} g(X(t), t) \right] dt + g(X(t), t) dW(t)$$

SDEs | Informal definition (cont.)

For the somewhat general case of N_x Itô stochastic differential equations, we have

$$\begin{aligned} dX_1(t) &= f_1(X_1(t), \dots, X_{N_x}(t), t) dt + g_1(X_1(t), \dots, X_{N_x}(t), t) dW_1(t) \\ &\dots = \dots \\ dX_{N_x}(t) &= f_{N_x}(X_1(t), \dots, X_{N_x}(t), t) dt + g_{N_x}(X_1(t), \dots, X_{N_x}(t), t) dW_{N_x}(t) \end{aligned}$$

Equivalently, in computational form

$$\begin{aligned} \Delta X_1(t) &= f_1(X_1(t), \dots, X_{N_x}(t), t) \Delta t + g_1(X_1(t), \dots, X_{N_x}(t), t) \underbrace{\sqrt{\Delta t} \xi_1(t)}_{\Delta W_1(t)} \\ &\dots = \dots \\ \Delta X_{N_x}(t) &= f_{N_x}(X_1(t), \dots, X_{N_x}(t), t) \Delta t + g_{N_x}(X_1(t), \dots, X_{N_x}(t), t) \underbrace{\sqrt{\Delta t} \xi_{N_x}(t)}_{\Delta W_{N_x}(t)} \end{aligned}$$

SDEs | Informal definition (cont.)

In matrix form, for $X(t) = (X_1(t), \dots, X_{N_x}(t))$ we have

$$\begin{bmatrix} \Delta X_1(t) \\ \vdots \\ \Delta X_{N_x}(t) \end{bmatrix} = \underbrace{\begin{bmatrix} f_1(X(t), t) \\ \vdots \\ f_{N_x}(X(t), t) \end{bmatrix}}_{N_x \times 1} \Delta t + \underbrace{\begin{bmatrix} g_1(X(t), t) \\ \vdots \\ g_{N_x}(X(t), t) \end{bmatrix}}_{N_x \times 1} \mathcal{N} \left(\underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \end{bmatrix}}_{N_x \times 1}, \underbrace{\begin{bmatrix} \Delta t & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & \Delta t \end{bmatrix}}_{N_x \times N_x} \right)$$

Note that the number of N_k of independent unit normals was constructed to be equal to the number N_x of process components, though nothing precludes the case $N_k \neq N_x$

SDEs | Informal definition (cont.)

For the more general case of N_x Itô equations and N_k Brownian motions, we have

$$dX_1(t) = f_1(X(t), t) dt + \sum_{n_k=1}^{N_k} g_{1,n_k}(X(t), t) dW_{n_k}(t)$$

... = ...

$$dX_{N_x}(t) = f_{N_x}(X(t), t) dt + \sum_{n_k=1}^{N_k} g_{N_x,n_k}(X(t), t) dW_{n_k}(t)$$

Equivalently, in computational form

$$\Delta X_1(t) = f_1(X(t), t) \Delta t + \sum_{n_k=1}^{N_k} g_{1,n_k}(X(t), t) \underbrace{\sqrt{\Delta t} \xi_1(t)}_{\Delta W_1(t)}$$

... = ...

$$\Delta X_{N_x}(t) = f_{N_x}(X(t), t) \Delta t + \sum_{n_k=1}^{N_k} g_{N_x,n_k}(X(t), t) \underbrace{\sqrt{\Delta t} \xi_{N_k}(t)}_{\Delta W_{N_k}(t)}$$

SDEs | Informal definition (cont.)

In matrix form, for $X(t) = (X_1(t), \dots, X_{N_x}(t))$ we have

$$\begin{bmatrix} \Delta X_1(t) \\ \vdots \\ \Delta X_{N_x}(t) \end{bmatrix} = \underbrace{\begin{bmatrix} f_1(X(t), t) \\ \vdots \\ f_{N_x}(X(t), t) \end{bmatrix}}_{N_x \times 1} \Delta t + \underbrace{\begin{bmatrix} g_{1,1}(X(t), t) & \cdots & g_{1,N_k}(X(t), t) \\ \vdots & \ddots & \vdots \\ g_{N_x,1}(X(t), t) & \cdots & g_{N_x,N_k}(X(t), t) \end{bmatrix}}_{N_x \times N_k} \mathcal{N} \left(\underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \end{bmatrix}}_{N_k \times 1}, \underbrace{\begin{bmatrix} \Delta t & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & \Delta t \end{bmatrix}}_{N_k \times N_k} \right)$$

That is,

$$dX(t) = f(X(t), t) dt + G(X(t), t) \mathcal{N}(0, dt I_{N_k})$$

SDEs | Informal definition (cont.)

$$dX_{n_x}(t) = f_{n_x}(X(t), t) dt + \sum_{n_k=1}^{N_k} g_{n_x, n_k}(X(t), t) dW_{n_k}(t)$$

The n_x -th Itô equation is the standard-form Langevin equation for the Markov process

The corresponding Stratonovich equation,

$$dX_{n_x}(t) = \left(f_{n_x}(X(t), t) + \frac{1}{2} \sum_{n_k=1}^{N_k} \sum_{n'_x=1}^{N_x} \frac{\partial g_{n_x, n_k}(X(t), t)}{\partial X^{n'_x}} g_{n_x, n_k}(X(t), t) \right) dt + \sum_{n_k=1}^{N_k} g_{n_x, n_k}(X(t), t) dW_{n_k}(t)$$

Stochastic differential equations

Examples

SDEs | Examples

Informal
definition

Examples

Fokker-Planck

Definition

Derivation

Example

$$dX(t) = f(X(t), t) dt + g(X(t), t) dW(t)$$

Let $f(X(t), t) = 0$ and $g(X(t), t) = 1$, we get the stochastic differential equation

$$X(t + dt) = X(t) + dW(t)$$

Using the informal treatment, we get the associated computational definition

$$X(t + \Delta t) = X(t) + \sqrt{\Delta t} \underbrace{\xi(t)}_{\mathcal{N}(0,1)}$$

We consider a time interval $\Delta t = 10^{-3}$ and an initial condition $X(0) = 0$

SDEs | Examples (cont.)

Informal
definition

Examples

Fokker-Planck

Definition

Derivation

Let $M_X(t)$ be the expected value of $X(t)$

$$M_X(t) = \langle X(t) \rangle$$

Let $V_X(t)$ be the variance of $X(t)$

$$\begin{aligned} V_X(t) &= \langle (X(t) - M_X(t))^2 \rangle \\ &= \langle X(t)^2 \rangle - M_X(t)^2 \end{aligned}$$

From $X(t + \Delta t) = X(t) + \sqrt{\Delta t}\xi(t)$, we have

$$\begin{aligned} M_X(t + \Delta t) &= \langle X(t) + \sqrt{\Delta t}\xi(t) \rangle \\ &= \langle X(t) \rangle + \langle \sqrt{\Delta t}\xi(t) \rangle \\ &= \langle X(t) \rangle + \sqrt{\Delta t} \underbrace{\langle \xi(t) \rangle}_{=0} \\ &= \langle X(t) \rangle \\ &= M_X(t) \end{aligned}$$

Because $X(0) = 0$, we have that $\langle X(0) \rangle = 0$ and thus $M_X(0) = 0$ and also $M_X(t) = 0$

SDEs | Examples (cont.)

From $X(t + \Delta t) = X(t) + \sqrt{\Delta t}\xi(t)$, we also have

$$\begin{aligned}
 V_X(t + \Delta t) &= \langle (X(t + \Delta t) - M_X(t + \Delta t))^2 \rangle \\
 &= \langle X(t + \Delta t)^2 \rangle - M_X(t + \Delta t)^2 \\
 &= \langle X(t + \Delta t)^2 \rangle - \underbrace{M_X(t)^2}_{=0} \\
 &= \langle X(t + \Delta t)^2 \rangle \\
 &= \langle (X(t) + \sqrt{\Delta t}\xi(t))^2 \rangle \\
 &= \langle X(t)^2 + 2X(t)\sqrt{\Delta t}\xi(t) + \Delta t\xi(t)^2 \rangle \\
 &= \langle X(t)^2 \rangle + 2\underbrace{\langle X(t)\sqrt{\Delta t}\xi(t) \rangle}_{=0} + \Delta t \underbrace{\langle \xi(t)^2 \rangle}_{=1} \\
 &= \langle X(t)^2 \rangle + \Delta t \\
 &= \langle X(t)^2 \rangle - M_X(t)^2 + \Delta t \\
 &= V_X(t) + \Delta t
 \end{aligned}$$

Because $X(0) = 0$, we have that $V_X(0) = 0$ and thus also $V_X(t) = 0$

SDEs | Examples (cont.)

We have that both expected value and variance of $(X(t))_{t \geq 0}$ do not depend on Δt

- It can be shown that higher-order moments are also independent of Δt

SDEs | Examples (cont.)

Example

Consider a diffusing particle whose position in time is $(X(t), Y(t), Z(t))$ for $t \geq 0$

We will show that the position at time $(t + dt)$ is related to position at time t

$$X(t + dt) = X(t) + (2D)^{1/2} dW_x(t)$$

$$Y(t + dt) = Y(t) + (2D)^{1/2} dW_y(t)$$

$$Z(t + dt) = Z(t) + (2D)^{1/2} dW_z(t)$$

Quantities $dW_x(t)$, $dW_y(t)$, and $dW_z(t)$ are increments of Brownian motions

Quantity $D = 10^{-4} \text{ mm}^2\text{sec}^{-1}$ is the diffusion constant

Using the informal treatment, we get the associated computational definitions

$$X(t + \Delta t) = X(t) + (2D)^{1/2} \sqrt{\Delta t} \xi_x(t)$$

$$Y(t + \Delta t) = Y(t) + (2D)^{1/2} \sqrt{\Delta t} \xi_y(t)$$

$$Z(t + \Delta t) = Z(t) + (2D)^{1/2} \sqrt{\Delta t} \xi_z(t)$$

We consider an interval $\Delta t = 10^{-2}$ and initial condition $(X(0), Y(0), Z(0)) = (0, 0, 0)$



SDEs | Examples (cont.)

Example

$$dX(t) = f(X(t), t) dt + g(X(t), t) dW(t)$$

Let $f(X(t), t) = 1$ and $g(X(t), t) = 1$, we get the stochastic differential equation

$$X(t + dt) = X(t) + dt + dW(t)$$

Using the informal treatment, we get the associated computational definition

$$X(t + \Delta t) = X(t) + \Delta t + \underbrace{\sqrt{\Delta t} \xi(t)}_{\mathcal{N}(0,1)}$$

We consider a time interval $\Delta t = 10^{-3}$ and an initial condition $X(0) = 0$

SDEs | Examples (cont.)

Let $M_X(t)$ be the expected value of $X(t)$

$$M_X(t) = \langle X(t) \rangle$$

Let $V_X(t)$ be the variance of $X(t)$

$$\begin{aligned} V_X(t) &= \langle (X(t) - M_X(t))^2 \rangle \\ &= \langle X(t)^2 \rangle - M_X(t)^2 \end{aligned}$$

From $X(t + \Delta t) = X(t) + \Delta t + \sqrt{\Delta t}\xi(t)$, we have

$$\begin{aligned} M_X(t + \Delta t) &= \langle X(t) + \Delta t + \sqrt{\Delta t}\xi(t) \rangle \\ &= \langle X(t) \rangle + \underbrace{\langle \Delta t \rangle}_{\Delta t} + \langle \sqrt{\Delta t}\xi(t) \rangle \\ &= \langle X(t) \rangle + \Delta t + \sqrt{\Delta t} \underbrace{\langle \xi(t) \rangle}_{=0} \\ &= \langle X(t) \rangle + \Delta t \\ &= M_X(t) + \Delta t \end{aligned}$$

Because $X(0) = 0$, we have that $\langle X(0) \rangle = 0$ and thus $M_X(0) = 0$ and also $M_X(t) = t$

SDEs | Examples (cont.)

From $X(t + \Delta t) = X(t) + \Delta t + \sqrt{\Delta t}\xi(t)$, we have

$$V_X(t) = t$$



The Fokker-Planck equation

Stochastic differential equations

Fokker-Planck equation

Consider a system that evolves according to an Itô stochastic differential equation

$$dX(t) = f(X(t), t) + g(X(t), t) dW(t)$$

We let the probability density function of the process be $p(x, t)$

The probability that $x \leq X(t) \leq x + dx$ is thus $p(x, t)dx$, at t

$$\int_{\Omega_x \equiv \mathbb{R}} p(x, t) dx = 1$$

We can determine $p(x, t)$ empirically, after computing a large number of realisations

- The fraction of realisations that arrived at $[x, x + \Delta x]$ at $t = 1$

It is possible to determine the equation of motion for the probability density $p(x, t)$

Fokker-Planck equation (cont.)

$$dX(t) = f(X(t), t) + g(X(t), t) dW(t)$$

It can be shown that density $p(x, t)$ evolves according to a partial differential equation

$$\frac{\partial p(x, t)}{\partial t} = \frac{\partial^2}{\partial x^2} \left(\frac{1}{2} g^2(x, t) p(x, t) \right) - \frac{\partial}{\partial x} (f(x, t) p(x, t))$$

This partial differential equation is the Fokker-Planck or Kolmogorov forward equation

It is useful to understand the PFE as a master equation for certain process $(X(t))_{t \geq 0}$

- Any Markov processes whose individual jumps are very small
- (Sample paths of $(X(t))_{t \geq 0}$ are continuous functions of t)

Fokker-Planck equation (cont.)

$$\frac{\partial p(x, t)}{\partial t} = \underbrace{\frac{\partial^2}{\partial x^2} \left(\frac{1}{2} g^2(x, t) p(x, t) \right)}_{\text{Diffusion}} - \underbrace{\frac{\partial}{\partial x} (f(x, t) p(x, t))}_{\text{Drift}}$$

The FP/KF equation is a convection–diffusion equation for the transfer of probability

- ↪ The first term has been called the ‘transport-’, ‘convection-’, or ‘drift-’ term
- ↪ The second term has been called the ‘diffusion-’ or ‘fluctuation-’ term

The FPE is a continuity equation for the probability density

$$\frac{\partial p(x, t)}{\partial t} = - \frac{\partial J(x, t)}{\partial x}$$

The Fickian probability flux $J(x, t)$

$$J(x, t) = f(x, t) p(x, t) - \frac{1}{2} \frac{\partial}{\partial x} (g(x, t) p(x, t))$$

Fokker-Planck equation (cont.)

$$\frac{\partial p(x, t)}{\partial t} = - \frac{\partial J(x, t)}{\partial x}$$

An equilibrium steady-state solution corresponds to the conditions

$$\begin{aligned} \frac{\partial p(x, t)}{\partial t} &= 0 \\ J(x, t) &\equiv 0 \end{aligned}$$

This leads to a first-order ODE for the equilibrium density $p_{\text{SS}}(x)$

$$f(x) p_{\text{SS}}(x) - \frac{1}{2} \frac{\partial}{\partial x} (g(x) p_{\text{SS}}(x)) = 0$$

Fokker-Planck equation (cont.)

$$dX(t) = f(X(t), t) dt + g(X(t), t) dW(t)$$

For the special case of $f(x, t) = 0$ and $g(x, t) = 1$, we get the differential equation

$$dX(t) = dW(t)$$

The associated Fokker-Planck equation is the diffusion equation

$$\frac{\partial p(x, t)}{\partial t} = \frac{1}{2} \frac{\partial^2 p(x, t)}{\partial x^2}$$

For the Fickian probability flux $J(x, t)$, we get

$$J(x, t) = - \frac{1}{2} \frac{\partial}{\partial x} (p(x, t))$$

Fokker-Planck equation (cont.)

At equilibrium, the steady-state solution for the probability density $p(x)$

$$p_{\text{SS}}(x) = C \frac{1}{g^2(x)} \exp \left[\int_0^x 2 \frac{f(x')}{g^2(x')} dx' \right] \quad (C > 0)$$

Constant C is whatever number makes p_{SS} integrate to one

$$C = \left(\int_{\Omega_x \equiv \mathbb{R}} \frac{1}{g^2(x)} \exp \left[\int_0^x 2 \frac{f(x')}{g^2(x')} dx' \right] \right)^{-1}$$

Fokker-Planck equation | Derivation

How to derive the equation of motion for the probability density function $p(x, t)$?

$$\frac{\partial p(x, t)}{\partial t} = \frac{\partial^2}{\partial x^2} \left(\frac{1}{2} g^2(x, t) p(x, t) \right) - \frac{\partial}{\partial x} (f(x, t) p(x, t))$$

We let $p(x, t|y, s)dx$ be the probability that $X(t) \in [x, x + dx]$, given $X(s) = y$

$$\dots\dots \underbrace{s \dots t} \dots\dots$$

Now, we want to consider a future time $t + \Delta t$, such that $s < t < t + \Delta t$

$$\dots\dots \underbrace{s \dots t} \underbrace{\dots t + \Delta t} \dots\dots$$

Given $X(s) = y$, the probability that $X(t + \Delta t) \in [z, z + \Delta z]$ is $p(z, t + \Delta t|y, s)dz$

Fokker-Planck equation | Derivation (cont.)

$$\dots\dots \underbrace{\quad y \quad}_{s} \dots \underbrace{\quad X(t) \quad}_{t} \dots \underbrace{\quad z \quad}_{t + \Delta t} \dots\dots$$

We are interested in the probability that X moves from y at time s , to z at $t + \Delta t$

We sum the probabilities of all possible paths through intermediate points $x(t)$

$$p(z, t + \Delta t | y, s) = \int_{\Omega_x} p(z, t + \Delta t | x, t) p(x, t | y, s) dx$$

To derive the Fokker-Planck equation, we consider a vanishingly small Δt

Fokker-Planck equation | Derivation (cont.)

$$p(z, t + \Delta t | y, s) = \int_{\Omega_x} p(z, t + \Delta t | x, t) p(x, t | y, s) dx$$

To proceed, we firstly multiply both sides by some smooth function $\varphi : \Omega_x \rightarrow \mathbb{R}$

$$p(z, t + \Delta t | y, s) \varphi(z) = \int_{\Omega_x} p(z, t + \Delta t | x, t) \varphi(z) p(x, t | y, s) dx$$

The integrating over Ω_x and rearranging, we get

$$\begin{aligned} \int_{\Omega_x} p(z, t + \Delta t | y, s) \varphi(z) dz &= \int_{\Omega_x} \int_{\Omega_x} p(z, t + \Delta t | x, t) \varphi(z) p(x, t | y, s) dx dz \\ &= \int_{\Omega_x} \left[\int_{\Omega_x} p(z, t + \Delta t | x, t) \varphi(z) dz \right] p(x, t | y, s) dx \end{aligned}$$

Fokker-Planck equation | Derivation (cont.)

Function $\varphi(\cdot)$ is chosen to be any arbitrary smooth function over the domain Ω_x

Thus, it has a Taylor's expansion about any point $z_0 \in \Omega_x$

$$\begin{aligned}\varphi(z) &= \sum_{k=0}^{\infty} \frac{1}{k!} \left. \frac{d\varphi(z)}{dz} \right|_{z=z_0} (z - z_0)^k \\ &= \varphi(z_0) + \varphi'(z_0)(z - z_0) + \frac{1}{2}\varphi''(z_0)(z - z_0)^2 + \dots\end{aligned}$$

We expand $\varphi(\cdot)$ about point $x \in \Omega_x$ and truncate, to get

$$\varphi(z) \approx \varphi(x) + \varphi'(x)(z - x) + \frac{1}{2}\varphi''(x)(z - x)^2$$

Fokker-Planck equation | Derivation (cont.)

Substituting, we get

$$\begin{aligned}& \int_{\Omega_x} p(x, y + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\int_{\Omega} p(z, t + \Delta t | x, t) \underbrace{\varphi(z)}_{\text{Taylor}} dz \right] p(x, t | y, s) dx \\ &= \int_{\Omega_x} \left[\int_{\Omega} p(z, t + \Delta t | x, t) \left(\varphi(x) + \varphi'(x)(z - x) + \varphi''(x) \frac{(z - x)^2}{2} \right) dz \right] p(x, t | y, s) dx\end{aligned}$$

Rearranging, we get

$$\begin{aligned}& \int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\int_{\Omega_x} \varphi(x) p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\int_{\Omega_x} \varphi'(x)(z - x) p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\int_{\Omega_x} \varphi''(x) \frac{(z - x)^2}{2} p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx\end{aligned}$$

Fokker-Planck equation | Derivation (cont.)

Further rearranging terms, we get

$$\begin{aligned} & \int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\varphi(x) \int_{\Omega_x} p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\varphi'(x) \int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) \int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \end{aligned}$$

We recognise the following identity,

$$\int_{\Omega_x} p(z, t + \Delta t | x, t) dz = 1$$

Fokker-Planck equation | Derivation (cont.)

$$\begin{aligned} & \int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\varphi(x) \underbrace{\int_{\Omega_x} p(z, t + \Delta t | x, t) dz}_{=1} \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\varphi'(x) \int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) \int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \end{aligned}$$

We also recognise another identity,

$$\int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz = \langle X(t + \Delta t) - \underbrace{x(t)}_z \rangle$$

We need the expected displacement,

$$X(t + \Delta t) - x(t) = \Delta x(t)$$

Fokker-Planck equation | Derivation (cont.)

Using the Itô stochastic differential equation, we get

$$\underbrace{X(t + \Delta t) - x(t)}_{\Delta x(t)} = f(x(t), t) \Delta t + g(x(t), t) \Delta W(t)$$

For the expected displacement, we get

$$\begin{aligned} \langle \Delta x(t) \rangle &= \langle f(x(t), t) \Delta t + g(x(t), t) \Delta W(t) \rangle \\ &= \langle f(x(t), t) \Delta t \rangle + \underbrace{\langle g(x(t), t) \Delta W(t) \rangle}_{\underbrace{\sqrt{\Delta t} \xi(t)}_{g(x(t), t) \times 0}} \\ &= f(x(t), t) \Delta t \end{aligned}$$

Thus, for the integral we have

$$\int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz = f(x(t), t) \Delta t$$

Fokker-Planck equation | Derivation (cont.)

$$\begin{aligned} &\int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\underbrace{\varphi(x) \int_{\Omega_x} p(z, t + \Delta t | x, t) dz}_{=1} \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\underbrace{\varphi'(x) \int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz}_{=f(x(t), t) \Delta t} \right] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) \int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz \right] p(x, t | y, s) dx \end{aligned}$$

We recognise one last identity for the squared displacement

$$\int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz = \left\langle \left(X(t + \Delta t) - \underbrace{x(t)}_z \right)^2 \right\rangle$$

Fokker-Planck equation | Derivation (cont.)

Using again the stochastic differential equation, we have

$$\underbrace{X(t + \Delta t) - x(t)}_{\Delta x(t)} = f(x(t), t) \Delta t + g(x(t), t) \Delta W(t)$$

For the expected squared displacement, we get

$$\begin{aligned} \langle (\Delta x(t))^2 \rangle &= \langle (f(x(t), t) \Delta t + g(x(t), t) \Delta W(t))^2 \rangle \\ &= \underbrace{\langle (f(x(t), t) \Delta t)^2 \rangle}_{f(x(t), t)^2 \Delta t^2} + \underbrace{\langle (2f(x(t), t) \Delta t) \left(g(x(t), t) \underbrace{\Delta W(t)}_{\sqrt{\Delta t} \xi(t)} \right) \rangle}_{2f(x(t), t)g(x(t), t)\Delta t^{3/2} \times 1} \\ &\quad + \underbrace{\langle (g(x(t), t) \Delta W(t))^2 \rangle}_{g(x(t), t)^2 \times \Delta t \times 1} \\ &= g(x(t), t)^2 \Delta t + \mathcal{O}(\Delta t^2) \end{aligned}$$

Thus, for the integral we have

$$\int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz = g(x(t), t)^2 \Delta t$$

Fokker-Planck equation | Derivation (cont.)

$$\begin{aligned} &\int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx \\ &= \int_{\Omega_x} \left[\varphi(x) \underbrace{\int_{\Omega_x} p(z, t + \Delta t | x, t) dz}_{=1} \right] p(x, t | y, s) dx \\ &\quad + \int_{\Omega_x} \left[\varphi'(x) \underbrace{\int_{\Omega_x} (z - x) p(z, t + \Delta t | x, t) dz}_{=f(x(t), t) \Delta t} \right] p(x, t | y, s) dx \\ &\quad + \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) \underbrace{\int_{\Omega_x} (z - x)^2 p(z, t + \Delta t | x, t) dz}_{=g(x(t), t)^2 \Delta t} \right] p(x, t | y, s) dx \end{aligned}$$

Fokker-Planck equation | Derivation (cont.)

After substituting those integrals, we get

$$\begin{aligned} \int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx &= \int_{\Omega_x} [\varphi(x) 1] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} [\varphi'(x) f(x(t), t) \Delta t] p(x, t | y, s) dx \\ &+ \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) g(x(t), t)^2 \Delta t \right] p(x, t | y, s) dx \end{aligned}$$

After some manipulations, we get

$$\begin{aligned} \int_{\Omega_x} p(x, t + \Delta t | y, s) \varphi(x) dx &= \int_{\Omega_x} p(x, t | y, s) \varphi(x) dx \\ &+ \Delta t \int_{\Omega_x} \varphi'(x) f(x(t), t) p(x, t | y, s) dx \\ &+ \Delta t \int_{\Omega_x} \left[\frac{1}{2} \varphi''(x) g(x(t), t)^2 \right] p(x, t | y, s) dx \end{aligned}$$

Fokker-Planck equation | Derivation (cont.)

Reordering some terms, we get

$$\begin{aligned} \int_{\Omega_x} \frac{p(x, t + \Delta t | y, s) - p(x, t | y, s)}{\Delta t} \varphi(x) dx \\ = \int_{\Omega_x} \varphi'(x) f(x, t) p(x, t | y, s) dx + \int_{\Omega_x} \varphi''(x) \frac{g(x, t)^2}{2} p(x, t | y, s) dx \end{aligned}$$

Then, integrating by parts

$$\begin{aligned} \int_{\Omega_x} \frac{p(x, t + \Delta t | y, s) - p(x, t | y, s)}{\Delta t} \varphi(x) dx = \\ - \int_{\Omega_x} \varphi(x) \frac{\partial}{\partial x} (f(x, t) p(x, t | y, s)) dx + \int_{\Omega_x} \varphi(x) \frac{\partial^2}{\partial x^2} \left(\frac{g(x, t)^2}{2} p(x, t | y, s) \right) dx \end{aligned}$$

Fokker-Planck equation | Derivation (cont.)

Collecting terms and joining the integrals, we get

$$0 = \int_{\Omega_x} \varphi(x) \times \left\{ -\frac{p(x, t + \Delta t | y, s) - p(x, t | y, s)}{\Delta t} - \frac{\partial}{\partial x} (f(x, y)p(x, t | y, s)) + \frac{\partial^2}{\partial x^2} \left(\frac{g(x, t)^2}{2} p(x, t | y, s) \right) \right\} dx$$

As the function $\varphi(\cdot)$ is arbitrary (non-zero), the term within brackets must be zero

$$\frac{p(x, t + \Delta t | y, s) - p(x, t | y, s)}{\Delta t} = \frac{\partial^2}{\partial x^2} \left(\frac{g(x, t)^2}{2} p(x, t | y, s) \right) - \frac{\partial}{\partial x} (f(x, t)p(x, t | y, s))$$

Taking the limit $\Delta t \rightarrow 0$, we obtain the Fokker-Planck equation

$$\frac{\partial}{\partial t} p(x, t | y, s) = \frac{\partial^2}{\partial x^2} \left(\frac{g(x, t)^2}{2} p(x, t | y, s) \right) - \frac{\partial}{\partial x} (f(x, t)p(x, t | y, s))$$